AFRL-RW-EG-TP-2012-007



Mixture Model for Determination of Shock Equation of State

Jennifer L. Jordan¹ Melvin R. Baer²

¹Air Force Research Laboratory AFRL/RW, Eglin AFB, FL 32542

²Sandia National Laboratories
Albuquerque, New Mexico 87185-0836

25 July 2012

Interim Report

This paper was published in the Journal of Applied Physics, April 2012. One or more of the authors is a U.S. Government employee working within the scope of their position; therefore, the U.S. Government is joint owner of the work and has the right to copy, distribute, and use the work. Any other form of use is subject to copyright restrictions.

This work has been submitted for publication in the interest of the scientific and technical exchange. Publication of this report does not constitute approval or disapproval of the ideas or findings.

Distribution A: Approved for public release; distribution unlimited. Approval Confirmation 96 ABW/PA # 96ABW-2011-0401, dated August 26, 2011

AIR FORCE RESEARCH LABORATORY, MUNITIONS DIRECTORATE

Air Force Materiel Command ■ United States Air Force ■ Eglin Air Force Base

NOTICE AND SIGNATURE PAGE

Using Government drawings, specifications, or other data included in this document for any purpose other than Government procurement does not in any way obligate the U.S. Government. The fact that the Government formulated or supplied the drawings, specifications, or other data does not license the holder or any other person or corporation; or convey any rights or permission to manufacture, use, or sell any patented invention that may relate to them.

Qualified requestors may obtain copies of this report from the Defense Technical Information Center (DTIC) (http://www.dtic.mil).

AFRL-RW-EG-TP-2012-007 HAS BEEN REVIEWED AND IS APPROVED FOR PUBLICATION IN ACCORDANCE WITH ASSIGNED DISTRIBUTION STATEMENT.

FOR THE DIRECTOR:

//ORIGINAL SIGNED// //ORIGINAL SIGNED// //ORIGINAL SIGNED//

HOWARD G. WHITE, PhD CHRISTOPHER L. VARNER JENNIFER L. JORDAN, PhD

Technical Advisor Branch Chief Project Manager

Ordnance Division Energetic Materials Branch Energetic Materials Branch

This report is published in the interest of scientific and technical information exchange, and its publication does not constitute the Government's approval or disapproval of its ideas or findings.

This page intentionally left blank

REPORT DOCUMENTATION PAGE

Form Approved OMB No. 0704-0188

Public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing this collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden to Department of Defense, Washington Headquarters Services, Directorate for Information Operations and Reports (0704-0188), 1215 Jefferson Davis Highway, Suite 1204, Arlington, VA 22202-4302. Respondents should be aware that notwithstanding any other provision of law, no person shall be subject to any penalty for failing to comply with a collection of information if it does not display a currently valid OMB control number. PLEASE DO NOT RETURN YOUR FORM TO THE ABOVE ADDRESS.

4 DEDORT DATE (DD MM V)()()		2 DATEC COVERED (From To)
1. REPORT DATE (DD-MM-YYYY)	2. REPORT TYPE	3. DATES COVERED (From - To)
25 July 2012	Interim	January 2009 – August 2011
4. TITLE AND SUBTITLE		5a. CONTRACT NUMBER
Mixture Model for Determination	on of Shock Equation of State	5b. GRANT NUMBER
		5c. PROGRAM ELEMENT NUMBER 62102F
6. AUTHOR(S)		5d. PROJECT NUMBER
, ,	2	4347
Jennifer L. Jordan ¹ , Melvin R. Bae	er ²	5e. TASK NUMBER
		95
		5f. WORK UNIT NUMBER
		05
¹ Air Force Research Laboratory, A ² Sandia National Laboratories, Al		8. PERFORMING ORGANIZATION REPORT NUMBER
		AFRL-RW-EG-TP-2012-007
9. SPONSORING / MONITORING AGENC	CY NAME(S) AND ADDRESS(ES)	10. SPONSOR/MONITOR'S ACRONYM(S)
		AFRL-RW-EG
Air Force Research Laboratory, M	unitions Directorate	
Ordnance Division		11. SPONSOR/MONITOR'S REPORT
Energetic Materials Branch (AFRI	L/RWME)	NUMBER(S)
Eglin AFB FL 32542-5910	,	AFRL-RW-EG-TP-2012-007
Technical Advisor: Dr. Jennifer L.	Jordan	
40 DIOTDIDUTION / AVAIL ADULTY OTA		

12. DISTRIBUTION / AVAILABILITY STATEMENT

Distribution A: Approved for public release; distribution unlimited. Approval Confirmation 96 ABW/PA # 96ABW-2011-0401, Dated August 26, 2011

13. SUPPLEMENTARY NOTES

DISTRIBUTION STATEMENT INDICATING AUTHORIZED ACCESS IS ON THE COVER PAGE AND BLOCK 12 OF THIS FORM.

14. ABSTRACT

Simple mixture models for the prediction of shock equations of state (Hugoniot) are a necessary tool for characterization of multiple composites. A mixture model for determining the shock equation of state of composite materials is presented. The model is completely flexible allowing for multiple (>2) components. Additionally, error propagation analysis for the two component mixture model has been accomplished. The model predicts the equation of state to 5%-15% of the experimental data, which is comparable to variations realized in meso-scale modeling of similar materials.

15. SUBJECT TERMS

Mixture model, composite, particulate, polymer, shock, Hugoniot

16. SECURITY CLASSIFICATION OF:		17. LIMITATION OF ABSTRACT	18. NUMBER OF PAGES	19a. NAME OF RESPONSIBLE PERSON Jennifer L. Jordan	
a. REPORT	b. ABSTRACT	c. THIS PAGE	UL	16	19b. TELEPHONE NUMBER (include area code)
UNCLASSIFIED	UNCLASSIFIED	UNCLASSIFIED	OL	10	850-882-8992

This page intentionally left blank

Mixture model for determination of shock equation of state

Jennifer L. Jordan^{1,a)} and Melvin R. Baer²

Air Force Research Laboratory (AFRL/RWME), Eglin AFB, Florida 32542, USA

(Received 12 September 2011; accepted 9 March 2012; published online 19 April 2012)

Simple mixture models for the prediction of shock equations of state (Hugoniot) are a necessary tool for characterization of multiple composites. A mixture model for determining the shock equation of state of composite materials is presented. The model is completely flexible allowing for multiple (>2) components. Additionally, error propagation analysis for the two component mixture model has been accomplished. The model predicts the equation of state to 5%–15% of the experimental data, which is comparable to variations realized in meso-scale modeling of similar materials. [http://dx.doi.org/10.1063/1.3702873]

I. INTRODUCTION

As composites become increasingly used in applications where they are exposed to shock loading, a mixture model able to predict the shock equation of state based on the properties of the constituents is necessary. 1-13 However, models presented in the literature primarily account for only two component composites, 1-6 are restricted by geometry, 7 or require finite element software. 4-6,8 Bennett and Horie present a detailed mixture model for reactive powder composites, including calculating the volumetric mixture properties of the final compacted intermetallic material.¹⁴ Current trends in mesoscale modeling have resulted in predictive capability for the properties of composites, 15-17 which require either high quality microstructures for the materials or the ability to artificially generate such microstructures. Given the differences in materials, particularly for polymer-based particulate composites consisting of hard particles in soft binders, the preparation of samples for microstructural analysis can be difficult.

There is very little discussion in the literature about the properties of the individual constituents, and how well they are known, affect the predicted properties of the composites. Often the models used are so complex as to prohibit the propagation of error from the constituent input terms.

This paper presents a non-reactive mixture model for fully dense composites, which can be implemented in a mathematical code, such as Matlab. ¹⁸ The model is flexible in the number of components that can be accommodate and is non-specific in the types of materials, e.g., polymer matrix composites and metallic alloys can both be accommodated. The error in the constituent properties has been propagated through the two-constituent model to determine the error bars for the calculated equation of state. Several two- and three-constituent composites are investigated with the model and compared with experimental data. The paper builds upon previous work which investigated the equation of state of epoxy-based particulate composites. ¹⁹

II. MIXTURE MODEL

A. Derivation of mixture model

In determining the nonreactive EOS for a multiphase mixture, mixture theory has been successfully applied to a variety of multiphase materials. A mixture shock Hugoniot is defined by determining the relationship of shock velocity to particle velocity, U_s and u_p , respectively.

Using jump state and characteristics theory,¹ the shock velocity is, by definition, $U_s = c_s + u_p$ and across a shock the change in density, for a two constituent mixture, is constrained by

$$\frac{\phi_1^0 \rho_1^0 + \phi_2^0 \rho_2^0}{\phi_1 \rho_1 + \phi_2 \rho_2} = 1 - \frac{u_p}{U_s},\tag{1}$$

where ϕ_k (k=1,2) is the volume fraction, ρ_k (k=1,2) is the true material density, and the 0 superscript corresponds to the undisturbed state. In the limit of homogeneous flow for a two constituent mixture, the mixture sound speed, c_s , is given by 20

$$\frac{1}{c_s^2} = \frac{1}{(U_s - u_p)^2} = (\phi_1 \rho_1 + \phi_2 \rho_2) \left(\frac{\phi_1}{\rho_1 c_1^2} + \frac{\phi_2}{\rho_2 c_2^2} \right).$$
 (2)

For a saturated mixture,

$$\phi_1 + \phi_2 = 1 \tag{3}$$

and the true density for each constituent is determined using the shock Hugoniot for the pure constituents,

$$\frac{\rho_k^0}{\rho_k} = 1 - \frac{u_p}{U_{sk}} \tag{4}$$

and

$$U_{s,k} = c_{0,k} + s_1 u_{p,k} + s_2 u_{p,k}^2. (5)$$

Therefore, a mixture shock velocity (corresponding to a specified particle velocity) is determined by combining Eqs. (1) and (2) and the resulting transcendental relationship

²Sandia National Laboratories, Albuquerque, New Mexico 87185-0836, USA

a) Author to whom correspondence should be addressed. Electronic mail: Jennifer.jordan@eglin.af.mil.

(for a specific volume fraction) is solved using a standard root solving method. The mixture quantities for pressure, P, and density are then defined by

$$P = P^{0} + (\phi_{1}^{0}\rho_{1}^{0} + \phi_{2}^{0}\rho_{2}^{0})u_{p}U_{s}$$
 (6)

and Eq. (1).

For a three constituent mixture, similar relationships hold and the procedure is applied first treating two constituents as a new "1+2" constituent given that $\phi_1^0 > 0$; hence, one defines a mixture-averaged density and sound speed, ρ_{I2} and c_{I2} . A normalized volume fraction is given by

$$\alpha = \frac{\phi_1}{\phi_1 + \phi_2} \tag{7}$$

and

$$\rho_{12} = \alpha \rho_1 + (1 - \alpha)\rho_2 \tag{8}$$

the mixture relationships of "1+2" phases are similar to Eqs. (2) and (4) rewritten in terms of the normalized volume fraction,

$$\frac{\rho_{12}^0}{\rho_{12}} = 1 - \frac{u_p}{U_{s,12}} = \frac{c_{12}}{u_p + c_{12}} \tag{9}$$

and

$$\frac{1}{c_{12}^2} = \frac{1}{\left(U_{s,12} - u_p\right)^2} = \left(\alpha \rho_1 + (1 - \alpha)\rho_2\right) \left(\frac{\alpha}{\rho_1 c_1^2} + \frac{1 - \alpha}{\rho_2 c_2^2}\right). \tag{10}$$

Given the particle velocity, u_p , true densities for 1 and 2 are defined by their respective EOS in their pure state and α is solved combining Eqs. (9) and (10) and solving the resulting equation with a root solver.

Similar relationships are used for a new "1 + 3" constituent defining ρ_{I3} and c_{I3} by the following equations:

$$\beta = \frac{\phi_1}{\phi_1 + \phi_3},\tag{11}$$

$$\rho_{13} = \beta \rho_1 + (1 - \beta)\rho_3, \tag{12}$$

$$\frac{\rho_{13}^0}{\rho_{13}} = 1 - \frac{u_p}{U_{s,13}} = \frac{c_{13}}{u_p + c_{13}},\tag{13}$$

$$\frac{1}{c_{13}^2} = \frac{1}{(U_{s,13} - u_p)^2} = \left(\beta \rho_1 + (1 - \beta)\rho_3\right) \left(\frac{\beta}{\rho_1 c_1^2} + \frac{1 - \beta}{\rho_3 c_3^2}\right),\tag{14}$$

and β is solved combining Eqs. (13) and (14), and the resulting equation is solved using a root solver.

Finally, for a three constituent system, the saturation constraint requires

$$\phi_1 + \phi_2 + \phi_3 = 1 \tag{15}$$

and with Eqs. (7) and (11), the shocked volume fractions are defined for a specified u_p .

$$\phi_1 = \frac{\alpha\beta}{\beta(1-\alpha) + \alpha},\tag{16a}$$

$$\phi_2 = \frac{\beta(1-\alpha)}{\beta(1-\alpha) + \alpha},\tag{16b}$$

and

$$\phi_3 = \frac{\alpha(1-\beta)}{\beta(1-\alpha) + \alpha}.$$
 (16c)

Conservation of mass across the jump state (shock) requires

$$\frac{\phi_1^0 \rho_1^0 + \phi_2^0 \rho_2^0 + \phi_3^0 \rho_3^0}{\phi_1 \rho_1 + \phi_2 \rho_2 + \phi_3 \rho_3} = 1 - \frac{u_p}{U_{s,123}}.$$
 (17)

Hence, the wave speed for a three phase mixture, $U_{s,123}$, is defined as a function of particle velocity, u_p , producing the shock Hugoniot for a three phase system.

A similar procedure is applied to a multi-component mixture containing four or more constituents. In this general procedure, a reference component, subscripted by "1" such that $\phi_1^0 > 0$, and normalized variables are defined,

$$\alpha_j = \frac{\phi_j}{\phi_1 + \phi_j},\tag{18}$$

$$\rho_{1i} = \alpha_i \rho_i + (1 - \alpha_i) \rho_1, \tag{19}$$

$$\frac{\rho_{1j}^0}{\rho_{1j}} = 1 - \frac{u_p}{U_{s,1j}} = \frac{c_{1j}}{u_p + c_{1j}},\tag{20}$$

and

$$\frac{1}{c_{1j}^2} = \frac{1}{(U_{s,1j} - u_p)^2} = \left(\alpha_j \rho_j + (1 - \alpha_j) \rho_1\right) \left(\frac{\alpha_j}{\rho_j c_j^2} + \frac{1 - \alpha_j}{\rho_1 c_1^2}\right). \tag{21}$$

The solution for all α_j (where $\alpha_j < 1$) can be determined upon specifying equations of state for each individual constituent and solving the sets of transcendental Equations (18)–(21). Since saturation implies that $\sum_{k=1}^{N} \phi_k = 1$, then

$$\phi_k = \frac{\alpha_k \phi_1}{1 - \alpha_k}, \quad \text{for } k = 2, N$$
 (22)

hence,

$$\phi_1 = \frac{1}{\left(1 + \sum_{k=2}^{N} \frac{\alpha_k}{1 - \alpha_k}\right)} \tag{23a}$$

and

$$\phi_j = \frac{\alpha_j}{(1 - \alpha_j) \left(1 + \sum_{k=2}^N \frac{\alpha_k}{1 - \alpha_k}\right)}$$
(23b)

and the overall density change is then related to the mixture wave speed to define a mixture equation of state (i.e., shock Hugoniot) according to

$$\frac{\sum_{k=1}^{N} \phi_k^0 \rho_k^0}{\sum_{k=1}^{N} \phi_k \rho_k} = 1 - \frac{u_p}{U_{s,mix}}.$$
 (24)

B. Uncertainty analysis in mixture relationships

Along the Hugoniot, the mixture pressure, P, is related to the average particle velocity, u_p , and the mixture wave speed, U_s , by

$$P = \rho_0^{mix} u_n U_s, \tag{25}$$

where ρ_0^{mix} is the undisturbed density. At a fixed particle velocity, it is readily seen that

$$\frac{\Delta P}{P} = \frac{\Delta U_s}{U_c} \tag{26}$$

or the uncertainty associated with estimating pressure is directly related to the uncertainty in the mixture wave speed. Hence, the variations of mixture wave speed are assessed with regard to uncertainties of the individual constituents of the mixture.

For a two component mixture, consider

$$U_s = f(\rho_1^0, \rho_2^0, \phi_1^0, \phi_1, U_{s1}, U_{s2})$$
(27)

at conditions of fixed u_p . Following Ref. 21, the appropriate relationship for the uncertainty of mixture wave speed, dU_s , in terms of the uncertainties of variables in Eq. (27) is

$$dU_{s} = \left[\left(\frac{\partial U_{s}}{\partial \rho_{1}^{0}} d\rho_{1}^{0} \right)^{2} + \left(\frac{\partial U_{s}}{\partial \rho_{2}^{0}} d\rho_{2}^{0} \right)^{2} + \left(\frac{\partial U_{s}}{\partial \phi_{1}^{0}} d\phi_{1}^{0} \right)^{2} + \left(\frac{\partial U_{s}}{\partial \phi_{1}} d\phi_{1} \right)^{2} + \left(\frac{\partial U_{s}}{\partial U_{s1}} dU_{s1} \right)^{2} + \left(\frac{\partial U_{s}}{\partial U_{s2}} dU_{s2} \right)^{2} \right]^{1/2}.$$

$$(28)$$

The function for U_s is defined by the relationship for the mixture wave speed, i.e., Eq. (2), which, knowing that

$$\frac{\rho}{\rho_0} = \frac{U_s}{U_s - u_p},\tag{29}$$

can be rearranged as

$$\frac{1}{U_s(U_s - u_p)} = \rho_0^{mix} \left\{ \frac{\phi_1}{\rho_1^0 U_{s1}(U_{s1} - u_p)} + \frac{1 - \phi_1}{\rho_2^0 U_{s2}(U_{s2} - u_p)} \right\}.$$
(30)

Partial differentiation of Eq. (30) in terms of each of the variables in Eq. (28) yields

$$dU_s = \left\{ (a_1 d\rho_1^0)^2 + (a_2 d\rho_2^0)^2 + (a_3 dU_{s1})^2 + (a_4 dU_{s2})^2 + (a_2 d\phi_1^0)^2 + (a_6 d\phi_1)^2 \right\}^{1/2}, \quad (31a)$$

where

$$a_{1} = \frac{(U_{s}^{2} - U_{s}u_{p})^{2}}{u_{p} - 2U_{s}} \left(\frac{-(1 - \phi_{1}^{0})\rho_{2}^{0}\phi_{1}}{(\rho_{1}^{0})^{2}U_{s1}(U_{s1} - u_{p})} + \frac{\phi_{1}^{0}(1 - \phi_{1})}{\rho_{2}^{0}U_{s2}(U_{s2} - u_{p})} \right),$$
(31b)

$$a_{2} = \frac{(U_{s}^{2} - U_{s}u_{p})^{2}}{u_{p} - 2U_{s}} \left(\frac{(1 - \phi_{1}^{0})\phi_{1}}{\rho_{1}^{0}U_{s1}(U_{s1} - u_{p})} - \frac{\phi_{1}^{0}\rho_{1}^{0}(1 - \phi_{1})}{(\rho_{2}^{0})^{2}U_{s2}(U_{s2} - u_{p})} \right),$$
(31c)

$$a_3 = \frac{(U_s^2 - U_s u_p)^2}{u_p - 2U_s} \left(\frac{\phi_1 \rho_0}{\rho_1^0}\right) \left(\frac{2U_{s1} - u_p}{(U_{s1}^2 - U_{s1} u_p)^2}\right), \quad (31d)$$

$$a_4 = \frac{(U_s^2 - U_s u_p)^2}{u_p - 2U_s} \left(\frac{(1 - \phi_1)\rho_0}{\rho_2^0} \right) \left(\frac{2U_{s2} - u_p}{(U_{s2}^2 - U_{s2} u_p)^2} \right), \tag{31e}$$

$$a_{5} = \frac{(U_{s}^{2} - U_{s}u_{p})^{2}}{u_{p} - 2U_{s}} (\rho_{1}^{0} - \rho_{2}^{0}) \times \left(\frac{\phi_{1}}{\rho_{1}^{0}U_{s1}(U_{s1} - u_{p})} + \frac{1 - \phi_{1}}{\rho_{2}^{0}U_{s2}(U_{s2} - u_{p})}\right), \quad (31f)$$

$$a_{6} = \frac{(U_{s}^{2} - U_{s}u_{p})^{2}}{u_{p} - 2U_{s}} \rho_{0} \left(\frac{1}{\rho_{1}^{0}U_{s1}(U_{s1} - u_{p})} - \frac{1}{\rho_{2}^{0}U_{s2}(U_{s2} - u_{p})} \right).$$
(31g)

All of the terms in Eq. (31a) are known, except for the error in the volume fraction at pressure, $d\phi_1$, since the volume fraction at pressure is a derived quantity. However, this error can be determined by differentiating Eq. (1),

$$d\phi_1 = \left\{ (b_1 dU_s)^2 + (b_2 d\phi_1^0)^2 + (b_3 d\rho_1^0)^2 + (b_4 d\rho_2^0)^2 + (b_5 dU_{s1})^2 + (b_6 dU_{s2})^2 \right\}^{1/2}.$$
 (32a)

where

$$b_1 = \frac{-u_p}{(U_s - u_p)^2 \delta_{p12}},\tag{32b}$$

$$b_2 = \frac{(\rho_1^0 - \rho_2^0)U_s}{(U_s - u_p)\rho_0 \delta_{p12}},$$
 (32c)

$$b_{3} = \frac{\left(\frac{\rho_{0}U_{s}}{U_{s} - u_{p}} - \frac{\rho_{2}^{0}U_{s2}}{U_{s2} - u_{p}}\right)\left(\frac{U_{s1}}{(U_{s1} - u_{p})}\right)}{-(\rho_{1}^{0})^{2}},$$
 (32d)

$$b_4 = \frac{\left(\frac{\rho_0 U_s}{(U_s - u_p)} - \frac{\rho_2^0 U_{s1}}{(U_{s1} - u_p)}\right) \left(\frac{U_{s2}}{(U_{s2} - u_p)}\right)}{-\left(\delta_{p12}\rho_0\right)^2}, \quad (32e)$$

$$b_5 = \left(\frac{-\rho_1^0 u_p}{(U_{s1} - u_p)^2}\right) \frac{\left(\frac{\rho_0 U_s}{(U_s - u_p)} - \frac{\rho_2^0 U_{s2}}{(U_{s2} - u_p)}\right)}{-(\delta_{p12}\rho_0)^2}, \quad (32f)$$

TABLE I. Constituent and composite parameters for uranium with 13.4 wt. % rhodium.

Material	$\rho_0 \\ (\text{g/cm}^3)$	C ₀ (km/s)	S	ϕ^0	$\mathrm{d} ho_0$	$d\phi^0$	dU_{s}	Refs.
U	18.93	2.51	1.51	0.735	0.05	0.001	0.061	22
Rh	12.429	4.79	1.36	0.265	0.05		0.058	22
U-13.4 wt. % Rh	17.20	2.74	1.49					22

$$b_6 = \left(\frac{-\rho_2^0 u_p}{(U_{s2} - u_p)^2}\right) \frac{\left(\frac{\rho_0 U_s}{(U_s - u_p)} - \frac{\rho_1^0 U_{s1}}{(U_{s1} - u_p)}\right)}{-(\delta_{p12}\rho_0)^2}, \quad (32g)$$

and

$$\delta_{p12} = \frac{1}{\rho_0} \left[\frac{\rho_1^0 U_{s1}}{(U_{s1} - u_p)} - \frac{\rho_2^0 U_{s2}}{(U_{s2} - u_p)} \right]. \tag{33}$$

Equation (30) can be substituted into Eq. (31a) and rearranged to solve for dU_s ,

$$dU_{s} = \left\{ \left(\frac{1}{1 - a_{6}^{2}b_{1}^{2}} \right) \left[(a_{1}^{2} + a_{6}^{2}b_{3}^{2})(d\rho_{1}^{0})^{2} + (a_{2}^{2} + a_{6}^{2}b_{4}^{2})(d\rho_{2}^{0})^{2} + (a_{3}^{2} + a_{6}^{2}b_{55}^{2})(dU_{s1})^{2} + (a_{4}^{2} + a_{6}^{2}b_{6}^{2})(dU_{s2})^{2} + (a_{5}^{2} + a_{6}^{2}b_{2}^{2})(d\phi_{1}^{0})^{2} \right] \right\}^{1/2}.$$
(34)

The error propagation for systems containing three or more components is increasingly complex and is not considered in this paper.

III. COMPARISON WITH EXPERIMENTAL DATA

A. Two-constituent composites

1. Uranium-rhodium composite

Uranium-rhodium (U-13.4 wt. % Rh) was chosen as a model metal-metal composite. The properties for each constituent used in the model are given in Table I. The bulk sound speed, C_o , and S were determined from linear least squares fit to the available data as shown in Figs. 1(a) and 1(b) for uranium and rhodium, respectively. The deviation from these linear fits, dU_s , were also determined and are shown in Fig. 1 and Table I.

The predicted Hugoniot relationship for U-13.4 wt. % Rh was determined from the mixture model for two constituents presented above. It can be seen from Fig. 2(a) that the model slightly overpredicts the experimental data, with an average deviation, dU_s/U_s of 0.05, shown in Fig. 2(b). The linear fits for uranium and rhodium are shown for comparison.

The propagated error associated with the input parameters to the mixture model was determined from Eq. (34) and is shown as the dashed lines bracketing the solid line in Fig. 2(a). The error in density was determined from the differences in the reported density values for uranium, and the same error was used for the density in rhodium. The density of the composite determined from the mixture model calculation agrees with that reported for the composite. In order to determine the error in the input volume fraction, it was assumed that there was +/-1 unit error in the last reported significant figure. The error in the shock velocity for both constituents is the dU_s values shown in Figs. 1(c) and 1(d).

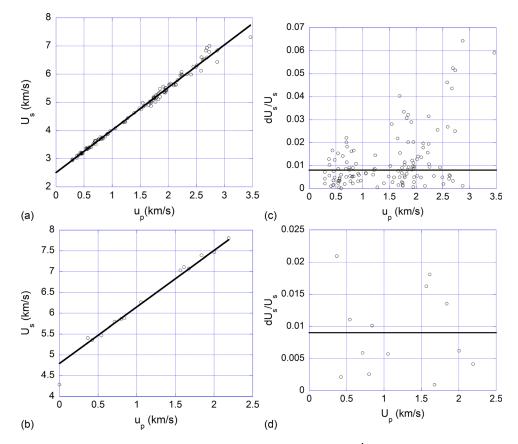


FIG. 1. Particle velocity vs shock velocity for (a) uranium and (b) rhodium, where the circles are the experimental data from Ref. 20 and the solid line is a linear fit to these experimental data, and dU_s/U_s for (c) uranium and (d) rhodium.

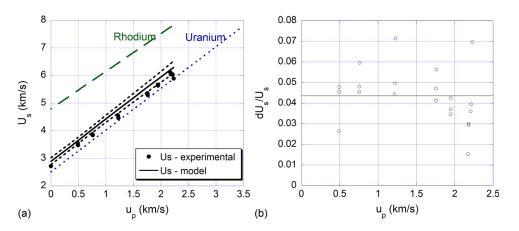


FIG. 2. (a) Particle velocity vs shock velocity and (b) dU_s/U_s for uranium, 13.4 wt.% rhodium experimental data (Ref. 20) compared with the mixture model calculation. The dashed lines on either side of the solid line in (a) indicated the error bars for the model calculation.

The experimental data²² seem to lie on or around the lower error bar for this composite.

2. Epoxy-aluminum and epoxy-alumina composites

Two epoxy-based two-phase particulate composites were considered: Al₂O₃-57 vol % epoxy^{23–26} and Al-55 vol % epoxy.²⁷ The former material has been extensively studied in the literature and the experimental data points used are combined from multiple references.^{23–26} The latter material had a limited number of experimental Hugoniot points available, but a comparison between a polymer-based composite containing a metal powder and one containing an oxide powder should be useful.

The properties of the constituent materials are presented in Table II (Refs. 22–32) and Fig. 3. The epoxy used in this study appears to undergo a phase change around $u_p = 2.84 \,\mathrm{km/s}$, Fig. 3(a). In order accommodate this phase change as well as a downturn in U_s at low values of u_p , a piecewise fit to the $U_s - u_p$ data for epoxy were used and is listed in Table III. In order to propagate the error from the input parameters through the model to determine the error bars in the output, the variations in the input parameters must be known. The variation in shock velocity, dU_s , was determined from the difference between the experimental data points and the linear least squares fit to the $U_s - u_p$ data. This variation is presented graphically in Figs. 3(d)–3(f) and

tabulated in Table II. Alumina deviates significantly from a linear U_s-u_p fit in the low pressure regime. Here the average dU_s , used in propagating the error for the mixture, may not be representative. The variation in density for each of the materials was determined from either the variation in reported densities in the literature or assuming a variation of +/-1 unit in the last reported significant figure.

The model prediction along with the experimental data for Al_2O_3 -57 vol % epoxy is presented in Fig. 4(a) and for Al-55 vol % epoxy in Fig. 4(c), along with the linear fits for the constituents. The model prediction underpredicts the data for Al_2O_3 -57 vol % epoxy by approximately 14% (Fig. 4(b)). The prediction for Al-55 vol % epoxy is better at approximately 5%; however, there exists limited data points available to compare to the model predictions.

B. Multi-component composites

1. 304 stainless steel

304 stainless steel was chosen as a representative three constituent composite, with the primary constituents, 74% Fe, 18% Ni, and 8% Cr,³³ considered in this study. The phase change in iron was neglected in the model, although it could be incorporated in the same method as the epoxy above. The properties for each constituent used in the model are given in Table IV. The bulk sound speed, C_o , and S were

TABLE II. Constituent and composite parameters for epoxy-aluminum and epoxy-alumina materials.

Material	$\rho_0 (\mathrm{g/cm}^3)$	C_0 (km/s)	S	ϕ_{epoxy}^0	$\mathrm{d} ho_0$	$d\phi^0$	$dU_s \\$	Ref.
Ероху	1.134–1.194 ^a				0.03		0.09	28–32
Aluminum	2.712	5.33	1.38		0.001		0.06	22
Alumina	3.976	8.14	1.28		0.007		0.24	22
Epoxy-Al ₂ O ₃ (300/100/20)	2.37	2.93	1.94	0.58		0.01		23
Epoxy-Al ₂ O ₃ (Al2)	2.377	2.88	1.99	0.57		0.01		24
Epoxy-Al ₂ O ₃ (Al3)	2.389			0.57		0.01		24
Epoxy-Al ₂ O ₃ (Al4)	2.391			0.57		0.01		24
Epoxy-Al ₂ O ₃ (fully loaded)	2.28	2.93	1.63	0.57		0.01		25
Epoxy-Al ₂ O ₃ (F1)	2.376	2.76	2.26			0.01		26
Epoxy-Al ₂ O ₃ (F2)	2.429	2.02	2.26			0.01		26
Epoxy-Al (Epoxy-45X81)	1.875			0.55		0.01		27

^aDependent on resin and curing agent.

10

TABLE III. Input parameters for mixture and mesoscale models.

Parameter	Epoxy Phase I	Epoxy Phase II	Epoxy Phase III	Epoxy Phase IV
u_p range	0-0.4	0.4-2.81	2.81-3.1	3.1-5.2
C_0 (km/s)	2.263	2.67	7.17	2.88
S	3.4704	1.55	0.01	1.35
S_2	-0.3947			

determined from linear least squares fits for Ni and Cr (Ref. 22) and used as reported in Ref. 34 for Fe.

The predicted Hugoniot relationship for 304 stainless steel was determined from the mixture model for three constituents presented above. It can be seen from Fig. 5 that the model slightly underpredicts the experimental data, with an average deviation, dU_s ,/ U_s of 0.05, shown in Fig. 5(b). The linear fits for iron, nickel, and chromium are shown for comparison. The model may be improved by the addition of the minor components in 304 stainless steel and inclusion of the

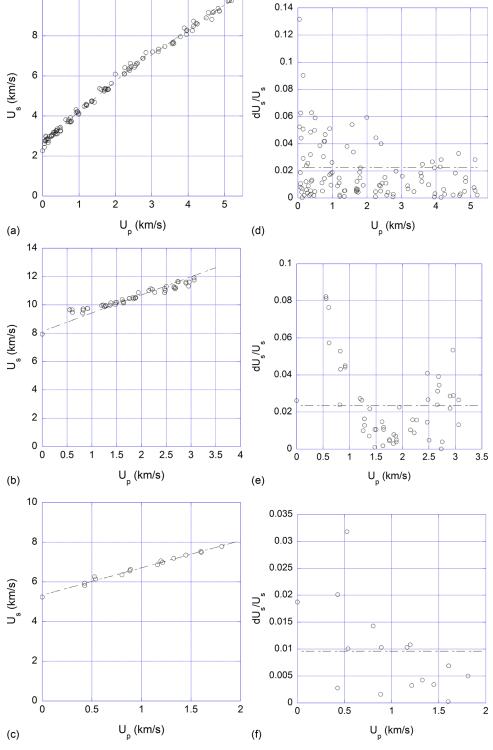


FIG. 3. Particle velocity vs shock velocity for (a) epoxy, (b) alumina, and (c) aluminum, where the circles are the experimental data from Refs. 20, 26–30, and the solid line is a linear least squares fit to this experimental data, and dU_s/U_s for (d) epoxy, (e) alumina, and (f) aluminum.

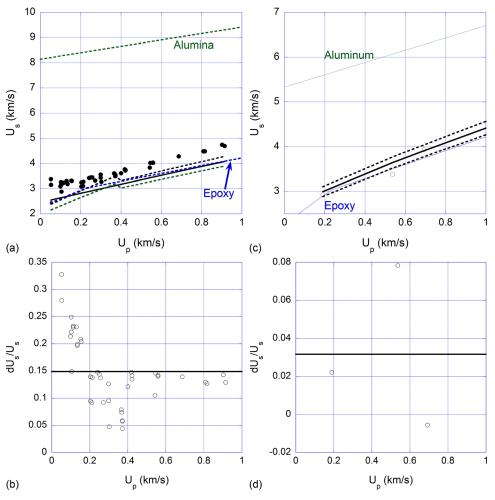


FIG. 4. (a) Particle velocity vs shock velocity and (b) dU_s/U_s for Al_2O_3 -57 vol % epoxy and (c) particle velocity vs shock velocity and (d) dU_s/U_s for Al-55 vol % epoxy experimental data compared with mixture model calculation. The dashed lines on either side of the solid line in (a) indicated the error bars for the model calculation.

phase change in iron; however, the agreement for this three component system is as good as that for the uranium-rhodium composite.

2. Epoxy-Al-Fe₂O₃ and epoxy-Al-MnO₂

Two multi-component epoxy-based composites have sufficient experimental data available in the literature, epoxy-Al-Fe₂O₃ (Ref. 35) and epoxy-Al-MnO₂.³⁶ The input data for epoxy and aluminum have been discussed in the previous section. The equation of state for MnO₂ was determined from that for pyrolusite, as presented in Ref. 22. All of the input data for the multi-phase composites is given in Table V. The equations of state for Fe₂O₃ and MnO₂ are presented graphically in Fig. 6 with their associated variation in shock wave speed. For both materials, epoxy was chosen as material "1," or the reference material, for Eqs. (7)–(15). However, the model is insensitive to the choice of the reference material, and simulations of the epoxy-Al-Fe₂O₃ system

TABLE IV. Constituent and composite parameters for 304 stainless steel and its constituents.

Material	$\rho_0 (\text{g/cm}^3)$	C ₀ (km/s)	S	ϕ^0	Refs.
Fe	7.85	3.57	1.92	0.74	34
Ni	8.875	4.59	1.44	0.18	22
Cr	7.119	5.2	1.43	0.08	22
304 SS	7.89	4.56	1.50		22

were conducted with each constituent as the reference material with no change in the predicted mixture behavior.

The experimental data and predicted model for epoxy-Al-Fe₂O₃ is shown in Fig. 7(a) along with the linear fits to the equations of state for the constituents. It can be seen that there is reasonably good agreement between the model and the experimental data. The \sim 7% error, Fig. 7(b), may be primarily due to the scatter in the experimental data rather than systematic differences between the experiments and the simulation. The model and experimental data are very close to the equation of state for epoxy, which is most likely due to the high volume fraction of epoxy in this system.

The predicted model and the experimental data for epoxy-Al-MnO₂ are presented in Fig. 7(c). For this composite, the density from the mixture model did not have good agreement with the actual density measured for the composite. The MnO₂ powders displayed porosity in Ref. 10, so the initial density of the MnO₂ constituent was modified by matching the composite mixture density to the experimentally measured composite density. The adjusted density is shown in parentheses in Table IV. With the adjusted density the agreement between the experimental data and the model is approximately 13%, which is in line with the agreement with two-phase epoxy based composites.

The experimental data in this system lies above the epoxy equation of state, rather than around it as in the epoxy-Al-Fe₂O₃ system. This indicates that the decreased volume fraction of epoxy in this system is reducing the effect of the

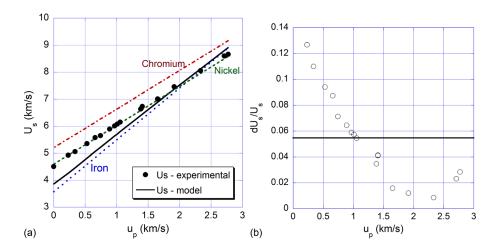


FIG. 5. (a) Particle velocity vs shock velocity and (b) dU_s/U_s for 304 stainless steel, experimental data (Ref. 20) compared with the mixture model calculation.

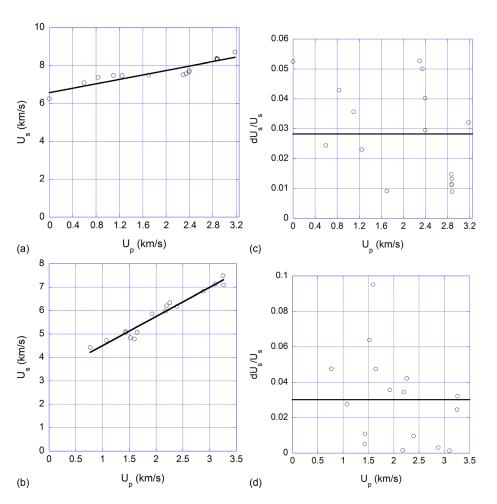


FIG. 6. Particle velocity vs shock velocity for (a) ${\rm Fe_2O_3}$ and (b) ${\rm MnO_2}$, where the circles are the experimental data from Ref. 20 and the solid line is a linear fit to these experimental data, and dU_s/U_s for (c) ${\rm Fe_2O_3}$ and (d) ${\rm MnO_2}$.

TABLE V. Constituent and composite parameters for epoxy-Al-Fe₂O₃ and epoxy-Al-MnO₂ materials.

Material	$\rho_0 (\mathrm{g/cm}^3)$	C ₀ (km/s)	S	ϕ^0_{epoxy}	$\mathrm{d} ho_0$	$d\phi^0$	dU_{s}	Refs.
Ероху	1.134–1.194 ^a				0.03		0.09	28–32
Aluminum	2.712	5.33	1.38		0.001		0.06	22
Iron oxide (Fe ₂ O ₃)	5.00	6.24	1.39		0.01		0.21	22
Manganese dioxide (MnO ₂)	5.03 (4.55)	3.261	1.24		0.01			22
Epoxy-Al-Fe ₂ O ₃	2.366			0.60		0.01		35
Epoxy-Al-MnO ₂	2.60	3.61 ^b	1.47 ^b	0.45	0.02	0.01		36

^aDependent on resin and curing agent.

^bLinear fit $0.7 \le u_p \le 1.7$ km/s.

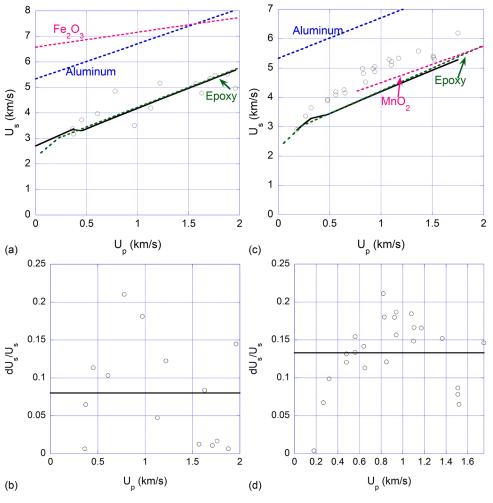


FIG. 7. (a) Particle velocity vs shock velocity and (b) dU_s/U_s for epoxy-Al-Fe₂O₃ and (c) particle velocity vs shock velocity and (d) dU_s/U_s for epoxy-Al-MnO₂ experimental (Refs. 30 and 31) data compared with the mixture model calculation. The dashed lines on either side of the solid line in (a) indicated the error bars for the model calculation.

epoxy binder on the composite behavior. The model does not reflect this same change in the influence of the epoxy binder. Increasing the aluminum volume fraction in the model to 0.40 (versus 0.24 actually in the composite), and decreasing the epoxy volume fraction by a comparable amount, results in a model that more accurately reflects the behavior of the composite. This may indicate that the volume fraction of the components varied in the samples or that the equation of state approximated for MnO₂ is not accurate.

IV. CONCLUSIONS

A general mixture model for determining the shock equation of state of composite materials is presented. The model is completely flexible allowing for multiple (>2) constituents in the composite. Additionally, error propagation analysis for the two component mixture model has been accomplished. For metallic composites, i.e., uraniumrhodium and 304 stainless steel, the average error between experiment and model is \sim 5%. The prediction error for epoxy-based composites increases to $\sim 10\%$ –15% based on the uncertainty in the epoxy equation of state, which dominates the material behavior due to its relatively large volume fraction. This mixture model allows for prediction of the equation of state of complex composite materials within 5%-15% of the experimental data, which is consistent with variations that are realized in meso-scale modeling of similar materials. The model can be implemented in MatLab and run in minutes on a desktop computer, which may provide an advantage over the longer running meso-scale simulations, which also require significant computational set-up to be able to run. This mixture model tool may be used to predict experiments involving complex materials that have not been previously tested, allowing for confidence in experimental timing and other parameters.

¹P. Embid and M. Baer, "Mathematical analysis of a two-phase model for reactive material," Continuum Mech. Thermodyn. **4**, 279 (1992).

²G. E. Duvall and S. M. Taylor, "Shock parameters in a two component mixture," J. Compos. Mater. **5**, 130–139 (1971).

E. Munson and K. W. Schuler, "Steady wave analysis of wave propagation in laminates and mechanical mixtures," J. Compos. Mater. 5, 286–304 (1971).
 M. R. Baer and J. W. Nunziato, "A two-phase mixture theory for the deflagration-to-detonation transition (ddt) in reactive granular materials," Int. J. Multiphase Flow 12, 861–889 (1986).

⁵A. K. Kapila, R. Menikoff, J. B. Bdzil, S. F. Son, and D. S. Stewart, "Two-phase modeling of deflagration-to-detonation transition in granular materials: Reduced equations," Phys. Fluids 13, 3002–3024 (2001).

⁶M. R. Baer, C. A. Hall, R. L. Gustavson, D. E. Hooks, and S. A. Sheffield, "Isentropic loading experiments of a plastic bonded explosive and constituents," J. Appl. Phys. **101**, 034906 (2007).

⁷F. K. Tsou and P. C. Chou, "Analytical study of Hugoniot in unidirectional fiber reinforced composites," J. Compos. Mater. 3, 500–514 (1969).

⁸L. M. Barker, "A model for stress wave propagation in composite materials," J. Compos. Mater. 5, 140–162 (1971).

⁹S. K. Garg and J. W. Kirsch, "Hugoniot analysis of composite materials," J. Compos. Mater. **5**, 428–445 (1971).

¹⁰S. K. Garg and J. W. Kirsch, "Steady shock waves in composite materials," J. Compos. Mater. 7, 277–285 (1973).

- ¹¹R. O. Davis and M. M. Cottrell, "Composite Hugoniot synthesis using the theory of mixtures," J. Compos. Mater. 5, 478–489 (1971).
- ¹²F. K. Tsou and P. C. Chou, "The control-volume approach to hugoniot of macroscopically homogeneous composites," J. Compos. Mater. 4, 526-537 (1970).
- ¹³J. E. Reaugh and E. L. Lee, "Shock Hugoniot behavior of mixed phases with widely varying shock impedences," Shock Compression of Condensed Matter -1997, edited by S. C. Schmidt, D. P. Dandekar, and J. W. Forbes
- ¹⁴L. S. Bennett and Y. Horie, "Shock-induced inorganic reactions and condensed phase detonations," Shock Waves 4, 127-136 (1994).
- ¹⁵R. A. Austin, D. L. McDowell, and D. J. Benson, "Numerical simulation of shock wave propagation in spatially-resolved particle systems," Modell. Simul. Mater. Sci. Eng. 14, 537–561 (2006).
- ¹⁶D. Eakins and N. N. Thadhani, "Discrete particle simulation of shock wave propagation in a binary Ni+Al powder mixture," J. Appl. Phys. 101, 043508 (2007).
- ¹⁷A. Fraser, J. P. Borg, J. L. Jordan, and G. Sutherland, "Exploring the micro-mechanical behavior of Al-MnO2-epoxy under shock loading while incorporating the epoxy phase transition," in Proceedings of the SEM Annual Conference, Albuquerque, NM, 2009.
- ¹⁸MATLAB (R2009b), The Mathworks, Inc., Natick, MA.
- ¹⁹J. L. Jordan, D. Dattelbaum, L. Ferranti, G. Sutherland, M. Baer, W. Richards, S. Sheffield, R. D. Dick, and N. N. Thadhani, "Shock equation of state of single constituent and multi-constituent epoxy-based particulate composites," Shock Compression of Condensed Matter -2009, edited by M. L. Elert, W. T. Buttler, M. D. Furnish, W. W. Anderson, and W. G. Proud(American Institute of Physics, 2009), pp. 1253–1256.
- ²⁰G. B. Wallis, One-Dimensional Two-Phase Flow (McGraw-Hill, 1969).
- ²¹H. W. Coleman and W. G. Steele, Jr., Experimentation and Uncertainty Analysis for Engineers (John Wiley & Sons, New York, 1989).
- ²²LASL Shock Hugoniot Data, edited by S. P. Marsh (University of California Press, Berkeley, 1980).
- ²³D. E. Munson, R. R. Boade, and K. W. Schuler, "Stress-wave propagation in Al2O3-epoxy mixtures," J. Appl. Phys. 49, 4797-4807 (1978).

- ²⁴R. E. Setchell, M. U. Anderson, and S. T. Montgomery, "Compositional effects on the shock-compression response of alumina-filled epoxy," J. Appl. Phys. 101, 083527 (2007).
- ²⁵J. C. F. Millett, N. K. Bourne, and D. Deas, "The equation of state of two alumina-filled epoxy resins," J. Phys. D: Appl. Phys. 38, 930–934 (2005).
- ²⁶R. E. Setchell and M. U. Anderson, J. Appl. Phys. **97**, 083518 (2005).
- ²⁷J. L. Jordan, E. B. Herbold, G. Sutherland, A. Fraser, J. Borg, and D. W. Richards, "Shock equation of state of multi-constituent epoxy-metal particulate composites," J. Appl. Phys. 109, 013531 (2011).
- ²⁸W. J. Carter and S. P. Marsh, "Hugoniot Equation of State of Polymer," Los Alamos National Laboratory Report LA-13006-MS, 1995.
- ²⁹D. E. Munson and R. P. May, "Dynamically determined high-pressure compressibilities of three epoxy resin systems," J. Appl. Phys. 43, 962-971 (1972).
- ³⁰M. U. Anderson, R. E. Setchell, and D. E. Cox, "Shock and release behavior of filled and unfilled epoxies," Shock Compression of Condensed Matter –1999, edited by M. D. Furnish, L. C. Chhabildas, and R. S. Hixson (American Institute of Physics, Melville, NY, 2000), pp. 551-554.
- ³¹J. C. F. Millett, N. K. Bourne, and N. R. Barnes, "The behavior of an epoxy resin under one-dimensional shock loading," J. Appl. Phys. **92**, 6590-6594- (2002).
- ³²P. J. Hazell, C. Stennett, and G. Cooper, "The shock and release behavior of an aerospace-grade cured aromatic amine epoxy resin," Polym. Compos. 29, 1106–1110 (2008).
- ³³J. R. Davis, *Metals Handbook* (ASM International, Materials Park, OH, 1998).
- ³⁴J. L. Jordan, L. Ferranti, R. A. Austin, R. D. Dick, J. R. Foley, N. N. Thadhani, D. L. McDowell, and D. J. Benson, "Equation of state of aluminumiron oxide-epoxy composite," J. Appl. Phys. 101, 093520 (2007).
- ³⁵M. A. Meyers, *Dynamic Behavior of Materials* (John Wiley and Sons, New York, 1994).
- ³⁶J. L. Jordan, D. M. Dattelbaum, G. Sutherland, D. W. Richards, S. A. Sheffield, and R. D. Dick, "Shock equation of state of a multi-phase epoxy-based composite (Al-MnO₂-epoxy)," J. Appl. Phys. 107, 103528 (2010).

DISTRIBUTION LIST AFRL-RW-EG-TP-2012-007

*Defense Technical Info Center 8725 John J. Kingman Rd Ste 0944 Fort Belvoir VA 22060-6218 AFRL/RWME (6) AFRL/RWOC-1 (STINFO Office)